# Ion non-Maxwellianity and its relation to turbulence and electrostatic waves

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[Valentini et al., 2016]



• We define the deviation from bi-Maxwellianity as:

$$\epsilon = \frac{1}{2n_i} \int_{v,\theta,\phi} |f_i(v,\theta,\phi) - f_1|$$
  
where

$$f_{\text{model}}(\mathbf{v}) = \frac{n_e}{\pi^{3/2} v_{i,\parallel}^3} \frac{T_{i,\parallel}}{T_{i,\perp}} \exp\left(-\frac{(v_{\parallel} - V_{\parallel})^2}{v_{i,\parallel}^2} - \frac{(v_{\perp,1} - V_{\perp})^2 + v_{\perp,2}^2}{v_{i,\parallel}^2 (T_{i,\perp}/T_{i,\parallel})}\right),$$

• We use ion moments to determine  $f_{model}$ :  $n = \int f(\mathbf{v}) d^3 v$ 

We use a bi-Maxwellian because we do not want a quantity that scales with  $T_{\parallel}/T_{\perp}$ .

#### Theory

 $f_{\text{model}}(v,\theta,\phi)|v^2\sin\theta dvd\theta d\phi,$ 

[Graham, et al., 2021b]

$$\mathbf{V} = \frac{1}{n} \int \mathbf{v} f(\mathbf{v}) d^3 v \qquad \mathbf{T} = \mathbf{P}/nk_B$$
$$\mathbf{P} = m \int (\mathbf{v} - \mathbf{V})(\mathbf{v} - \mathbf{V}) f(\mathbf{v}) d^3 v$$

#### $n = 20 \text{ cm}^{-3}$ T = 4 eV $T_{\perp}/T_{\parallel} = 0.7$ $\varepsilon$ = 0.12



#### **Example distributions**

 $n = 9 \text{ cm}^{-3}$ T = 5 eV $T_{\perp}/T_{\parallel} = 1.7$  $\varepsilon = 0.56$ 

### Solar wind example



 Example of ion Non-Maxwellianity in the solar wind.

$$\sqrt{Q} = \left(\frac{P_{12}^2 + P_{13}^2 + P_{23}^2}{P_{\perp}^2 + 2P_{\perp}P_{\parallel}}\right)^{1/2}$$

$$P_i = \begin{pmatrix} P_{\parallel} & P_{12} & P_{13} \\ P_{12} & P_{\perp} & P_{23} \\ P_{13} & P_{23} & P_{\perp} \end{pmatrix}$$
[Swisdak, 2016]

 Large *\varepsilon* occurs in turbulent regions in this example.

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#### Statistics



• Distribution of  $\varepsilon$  for all available data.  $\epsilon = \frac{1}{2n_i} \int_{\nu,\theta,\phi} ...$  $|f_i(v,\theta,\phi) - f_{\text{model}}(v,\theta,\phi)| v^2 \sin\theta dv d\theta d\phi,$ 

$$\sqrt{Q} = \left(\frac{P_{12}^2 + P_{13}^2 + P_{23}^2}{P_{\perp}^2 + 2P_{\perp}P_{\parallel}}\right)^{1/2}$$

Statistical relation between  $\varepsilon$  and sqrt(Q).

#### Temperature components

distributions by computing the eigenvalues/vectors of the temperature tensor.



[Servidio, et al., 2016]

• To further investigate the cause of high  $\varepsilon$ , we calculate 'proper' temperatures of the



### Example

log<sub>10</sub> B<sup>∕</sup>

nT<sup>2</sup> Hz<sup>-</sup>





 Cigar-shaped distributions are regularly oblique to **B**.

 Non-gyrotropic and non-Maxwellianity is primarily due to cigar-shaped distributions.

### Statistics



is associated with cigar-shaped distributions.

#### **Relation to turbulence**





 $|\Delta \mathbf{B}(t,\tau)|$  $\Delta \mathbf{B}(t,\tau) = \mathbf{B}(t+\tau) - \mathbf{B}(t)$ **PVI**  $\sqrt{\langle |\Delta \mathbf{B}(t,\tau)|^2 \rangle}$ 

[Greco, et al., 2008]





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#### Ion-acoustic waves in the solar wind



**Electron-ion** 

**Electron-electron-ion** 

**Electron-ion-ion** 

$$0 = 1 - \sum_{j} \frac{\omega_{pj}^2}{k^2 v_j^2} Z'\left(\zeta_j\right),$$

- Ion-acoustic waves can be generated by a variety of streaming instabilities.
- Based on estimated current densities, electron-ion-ion instability is most like to generate ion-acoustic waves.

[Graham, et al., 2021a]



#### Ion-acoustic wave examples

#### Non-Maxwellianity and ion-acoustic waves





 $\varepsilon = 0.5$ 

### Statistics



 $\epsilon$ LFR IA 8.0

• Slightly larger  $\varepsilon$  are observed in association with ion-acoustic waves.



## Conclusions

- Ion non-Maxwellianity is routinely observed in the solar wind.
- Ion non-Maxwellianity is not strongly correlation with local turbulent structures in the solar wind.
- Slightly enhanced non-Maxwellianity is observed at the same time as ion-acoustic waves. Kinetic instability of ion distributions is a plausible source of ion-acoustic waves.